Thermodynamics of Gambling Demons

Theory and Experiment

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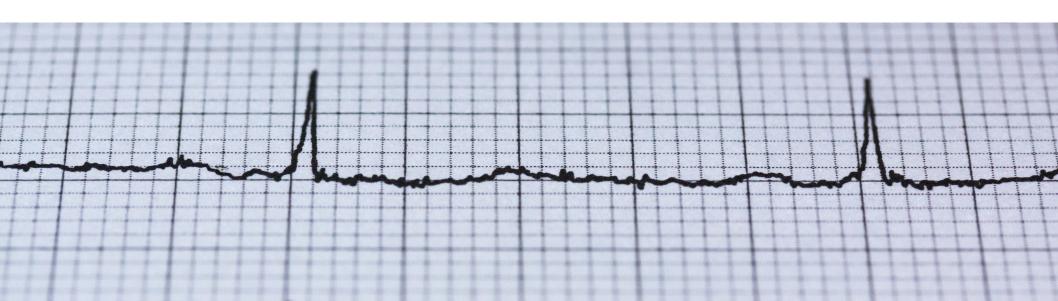






Outline of this talk:

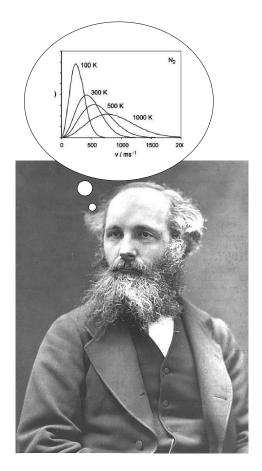
- Introduction
- Idea of gambling demons
- Martingale theory
- Theoretical results
- Experiment: single-electron box



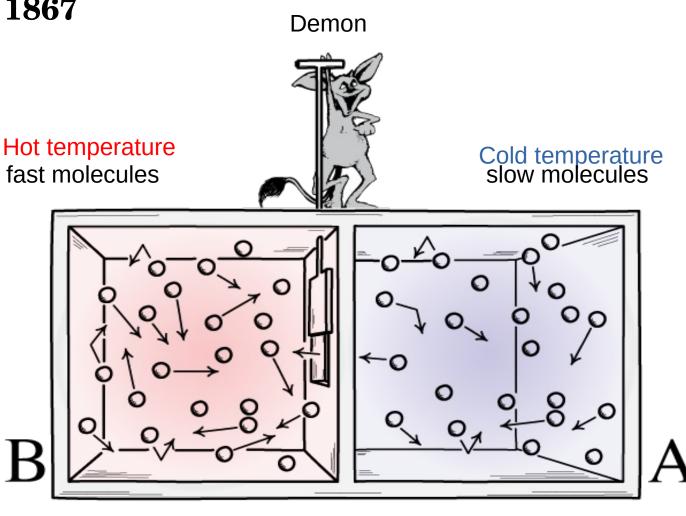






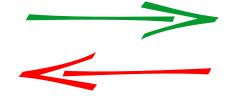


James C. Maxwell (1831 – 1879)



Spontaneaous heat flow

Paradoxical effect

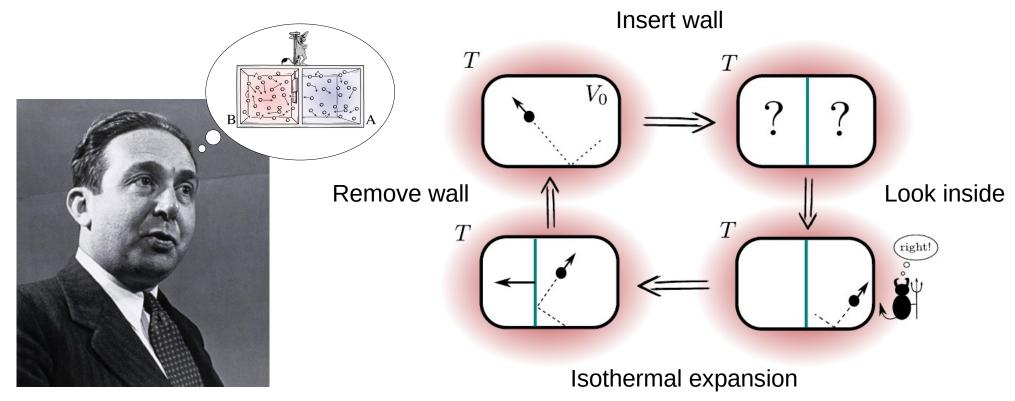


Drawing: Jonh D. Norton





Szilard's engine 1929



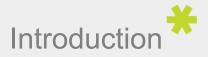
Leó Szilard (1898 - 1964)

Optimal work extraction per cycle:

$$W_{\text{ext}} = k_B T \log\left(\frac{V_0}{V_0/2}\right) = k_B T \log 2$$

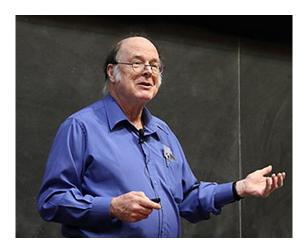
Szilard's conclusion: thermodynamic cost of measurement?





Informational Exorcism

Only apparent "violations" of second law. Take into account informational costs!



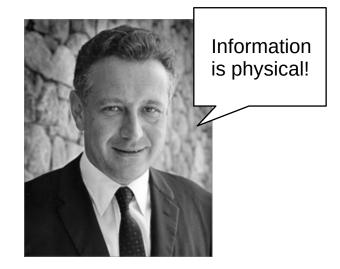
Charles H. Bennett

Landauer's Erasure 1961

 $W \ge k_B T \log 2$

Erasing information (1 bit of) has a minimum work cost even if doing it reversibly





Rolf W. Landauer (1927 - 1999)

Thermodynamics of feedback control:

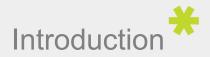
$$\Delta S - \frac{Q}{T} \ge -\mathcal{I}$$

information acquired /stored

Recall Jonh Bechoefer and Carlos Alvarez talks !

Theory Review: JMR Parrondo, JM Horowitz, and T Sagawa, Nat. Phys. (2015).

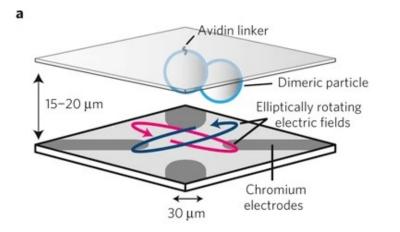




Experiments

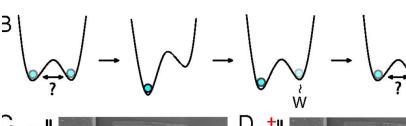
Colloidal particles in optical traps

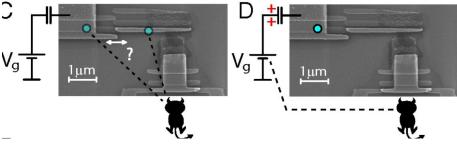
S Toyabe et al. Nat. Phys. 6 (2010)



Electronic devices

JV Koski et al. PNAS 111 (2014)



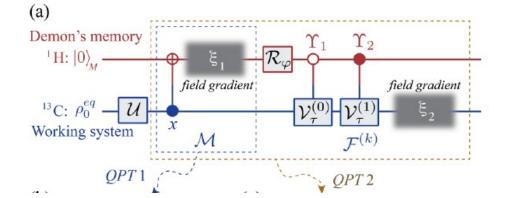


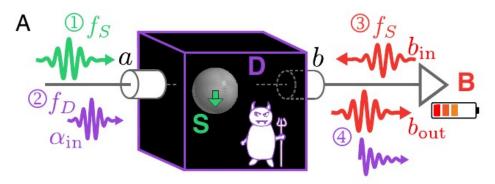
Nuclear spins and NMR spectrometry

PA Camati et al. PRL 117 (2016)

Circuit QED setups

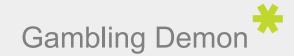
N Cottet et al. PNAS 114 (2017)





Experiments Review: S Ciliberto, E. Lutz, Physics Today (2015)





Gambling Demons

New version of Maxwell's demon based on stopping strategies:

(I) gather information about microscopic dynamics → ok

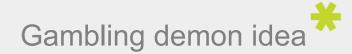
(II) feedback control (trapdoor / piston) —

Reduce to minimal expression! only stop (or not) the dynamics

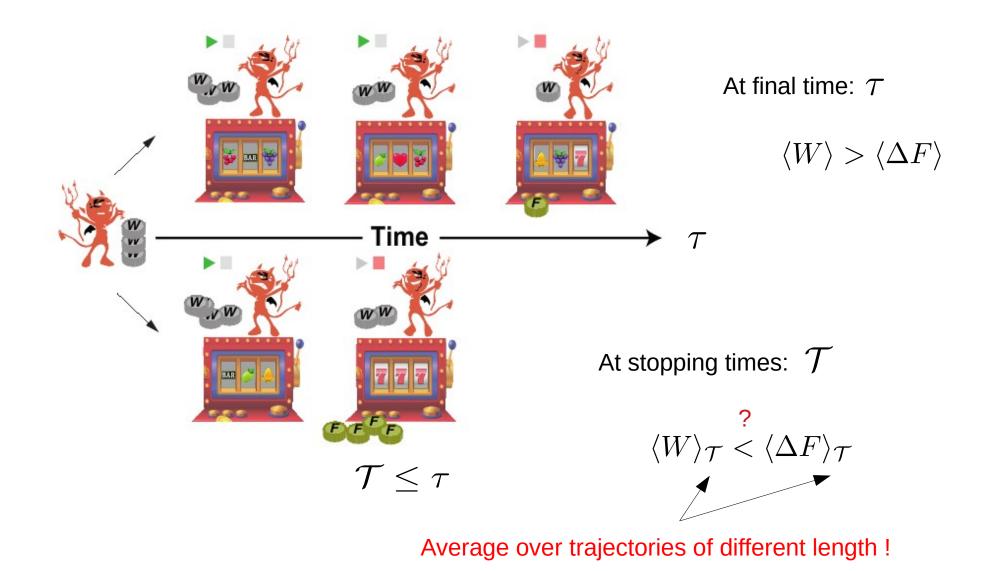


Players in the casino or agents in the market can decide to play or not a game, but cannot change the rules of the game (feedback).

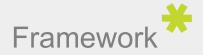




Analogous to a gambler in a slot machine:







Second law at the microscale: Fluctuation relations

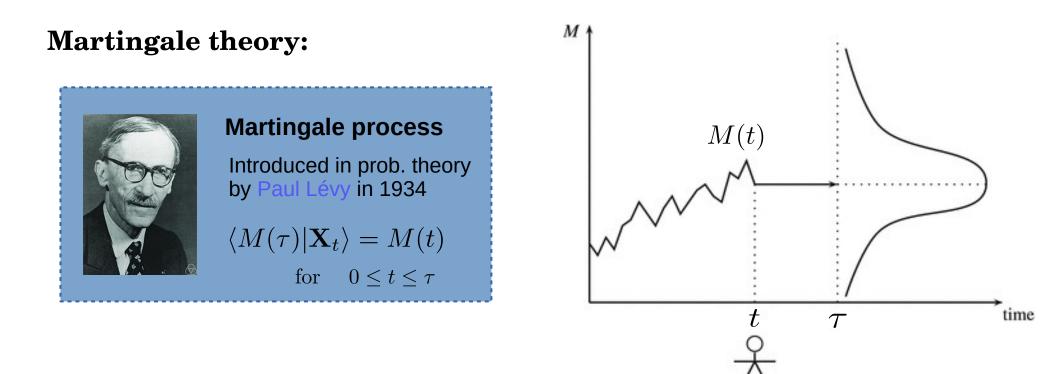
n = 1 Marrie System degree of freedom: X(t)_{det}(pA) 0 X(t)n = 0 follows stochastic paths $\mathbf{X}_{\tau} = \{X(t)\}_{t=0}^{\tau}$ -40 0.01 0.02 0.03 0.04 t(s) $S_{\text{tot}}(\tau) = \log\left(\frac{P(\mathbf{X}_{\tau})}{\tilde{P}(\mathbf{\tilde{X}}_{\tau})}\right) = \Delta S - \sum_{k} \frac{Q_{k}}{T_{k}}$ Stochastic entropy production: reservoirs system $S_{\rm tot} < 0$ $\langle e^{-S_{\text{tot}}} \rangle = 1 \qquad \langle S_{\text{tot}} \rangle \ge 0$ $P(S_{\rm tot})$ $S_{\rm tot} >$ exponential decay $S_{\rm tot}$

U. Seifert, Rep. Prog. Phys. (2012)

M. Esposito et al. Rev. Mod. Phys. (2009)





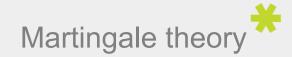


Martingales verify many interesting properties:

- Doob's optional stopping theorem: $\langle M(t) \rangle_{\mathcal{T}} = \langle M(0) \rangle$ at stochastic times \mathcal{T}
- Doob's maximal inequality: $\Pr[M_{\max}(\tau) \ge m] \le \langle M(\tau) \rangle / m$

maximum in an interval





Entropy production is an exponential martingale:

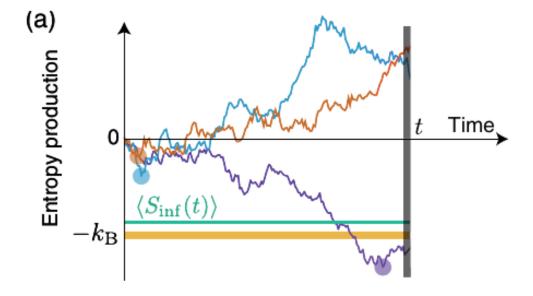
$$\langle e^{-S_{\rm tot}(\tau)} | \mathbf{X}_t \rangle = e^{-S_{\rm tot}(t)}$$

Nonequilibrium stationary processes

$$t = 0$$
 $\langle e^{-S_{\text{tot}}(\tau)} | X_0 \rangle = e^{-S_{\text{tot}}(0)} = 1$

- + Stopping-times fluctuation relation
 - $\langle e^{-S_{\text{tot}}(\mathcal{T})} \rangle = 1 \qquad \langle S_{\text{tot}}(\mathcal{T}) \rangle \ge 0$
- + Statistics of EP finite-time minimum

$$\Pr\left(S_{\min}(t) \le -\xi\right) \le e^{-\xi}$$



I. Neri, É. Roldán, F. Julicher PRX 7 (2017), R. Chétrite et al. EPL 124 (2019), I. Neri et al. JSM (2019) ...

if

Recent review: É.Roldán et al. arXiv:2210.09983 (2022)





Generalized stopping-times fluctuation relations

Extension of martingale theory for **driven systems** (with fixed protocol):

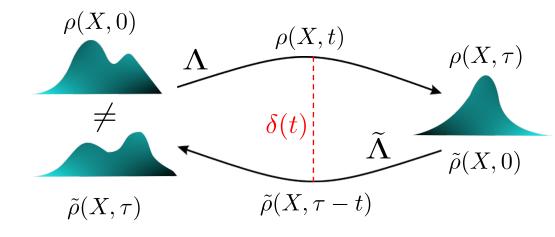
$$\langle e^{-S_{\text{tot}}(\tau)} | \mathbf{X}_t \rangle \neq e^{-S_{\text{tot}}(t)}$$
 but $\langle e^{-S_{\text{tot}}(\tau) - \delta(\tau)} | \mathbf{X}_t \rangle = e^{-S_{\text{tot}}(t) - \delta(t)}$

with the extra term
"stochastic asymmetry":
$$\delta(t) = \log\left(\frac{\rho[X(t), t]}{\tilde{\rho}[X(t), \tau - t]]}\right)$$

asymmetry of the prob. density under time-reversal

$$\begin{cases} \langle e^{-\beta[W-\Delta F]-\delta} \rangle_{\mathcal{T}} = 1 \end{cases}$$

Entropy production Extra term
$$\begin{cases} \langle W \rangle_{\mathcal{T}} - \langle \Delta F \rangle_{\mathcal{T}} \geq -k_B T \langle \delta \rangle_{\mathcal{T}} \end{cases}$$







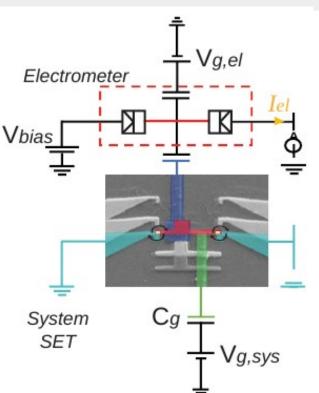
Single Electron Box (SEB): [PICO group (Helsinki)]

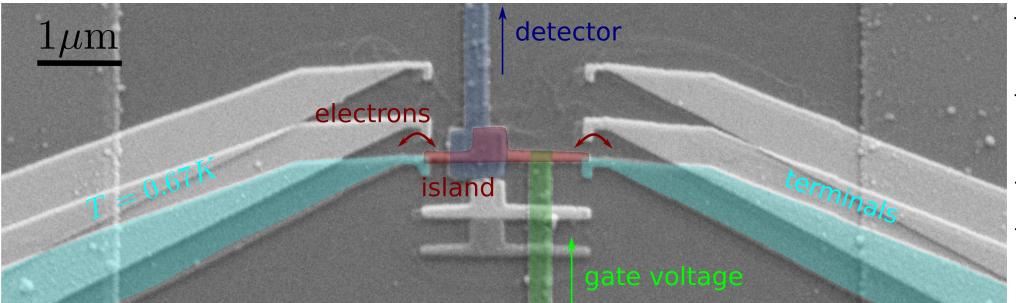
System: Cu island

Thermal reservoir: Al superconducting leads

Driving protocol: gate voltage following a linear ramp

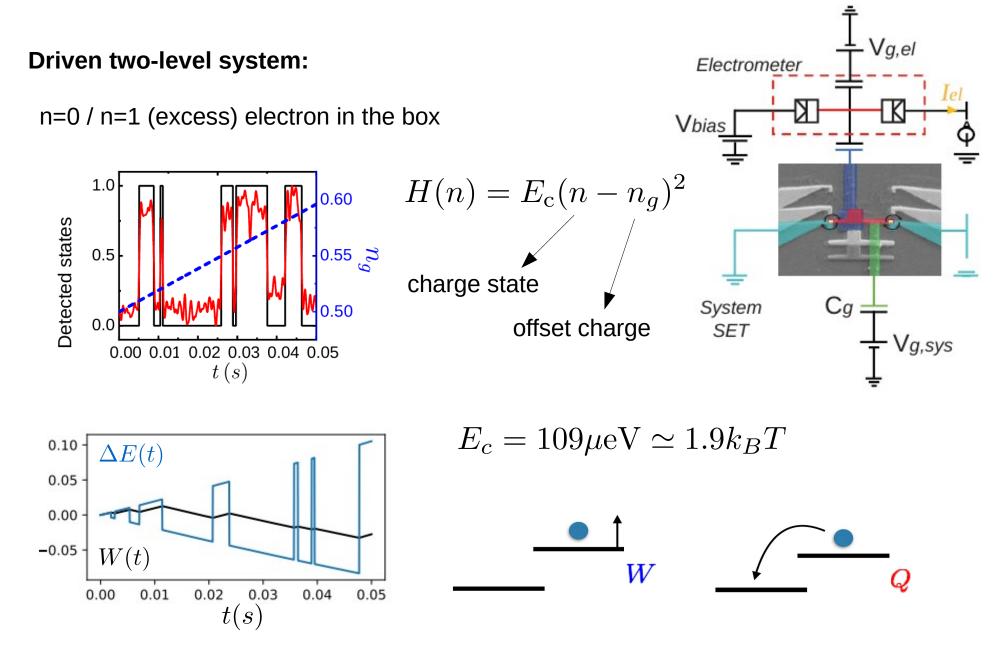
Detector: SET monitors tunneling events









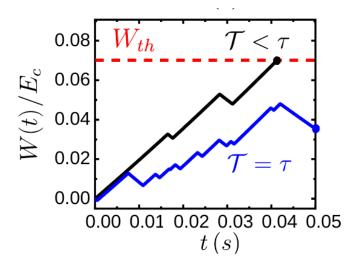


More details (setup): J.P. Pekola, Nat. Phys. **11** (2015) O. Maillet *et al*. PRL **122** (2019)

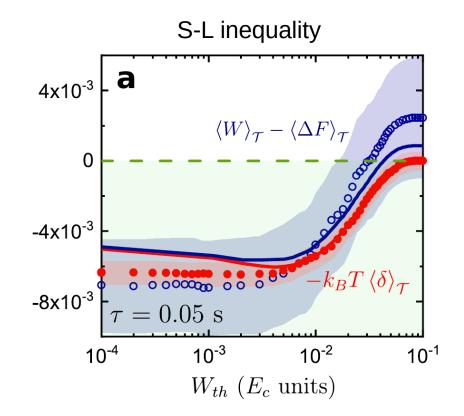


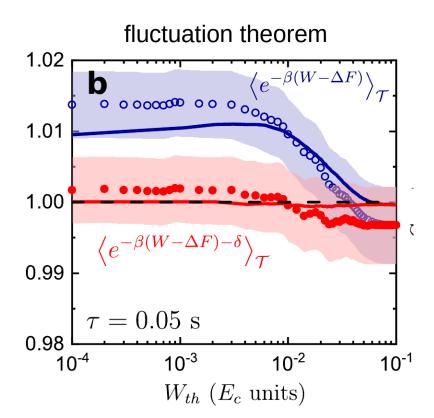
Strategy: Finite-horizon work first-passage times

Stop trajectories when the work reaches a threshold If threshold is not reached, we continue until the end.



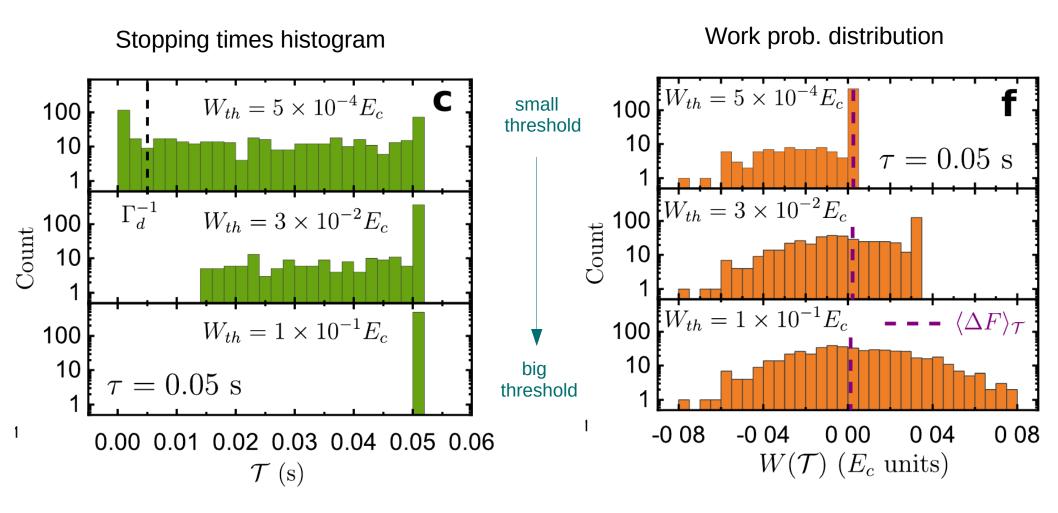
Experimental results:











We can extract work by stopping at small thresholds (on average) !

At large thresholds almost all trajectories stop at the final time: $~{\cal T}\simeq au$





Remarks:

 $\langle W \rangle_{\mathcal{T}} - \langle \Delta F \rangle_{\mathcal{T}} \geq -k_B T \langle \delta \rangle_{\mathcal{T}}$

Successful gambling when we have:

1. Time-reversal asymmetry

 $\langle \delta \rangle_{\mathcal{T}} > 0$

2. A good strategy / Information

monitor the system to decide when to stop

Outlook: compare with information inequalities ?

LETTERS https://doi.org/10.1038/s41567-019-0481-0

Large work extraction and the Landauer limit in a continuous Maxwell demon

M. Ribezzi-Crivellari 1,2 and F. Ritort 1,3*

nature

physics

 $\langle W \rangle - \langle \Delta F \rangle \ge -k_B T \mathcal{I}$

In the continuous limit:

 $\mathcal{T} \longrightarrow \infty$

Looser bounds

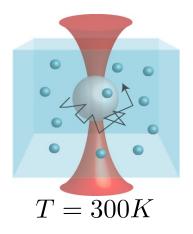


Other recent experiment:

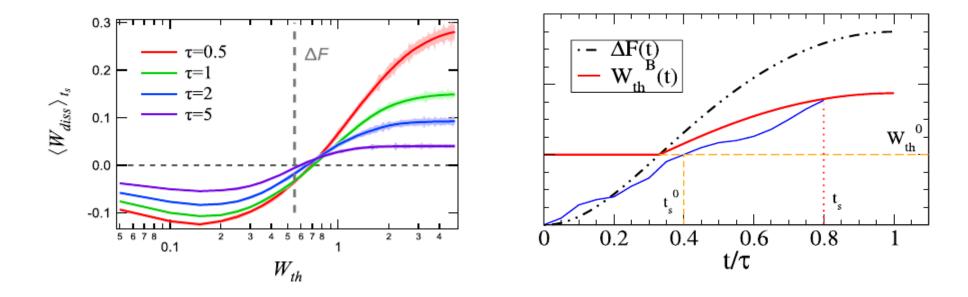
J.A.C. Albay, Y. Jun and P-Y Lai, PRR 5, 023115 (2023)

Overdamped Brownian particle under harmonic trapping (optical tweezers)

$$V(x,\lambda) = \frac{1}{2}\lambda(t)x^2$$
 $\lambda(t) = \lambda_0 [2 - \cos\left(\frac{\pi t}{\tau}\right)]$



Testing improvements modifying the stopping thresholds:







Main conclusions

- We introduced a "gambling demon" that stops a driven nonequilibrium thermodynamic process at stochastic times.
 - + Since it does not require feedback control, its applicability should be easier.
- We derived universal fluctuation relations for work and entropy production at stopping times and second-law-like inequalities.
 - + Valid in more general situations e.g. arbitrary strategies, several baths, etc.
- We tested our results experimentally in a single electron box, where we have found a simple and good strategy allowing work extraction.

+ Optimal strategies?











for your attention

For more information: Phys. Rev. Lett. 126, 080603 (2021)





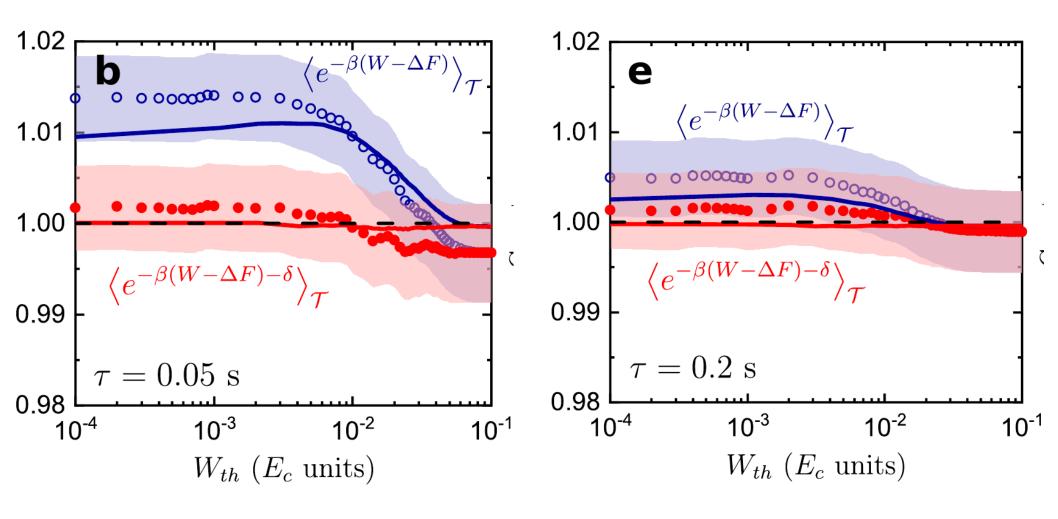








Experimental results: Fluctuation theorem at stopping times



Faster protocol (larger "violations")

Slower protocol (smaller "violations")





• Gambling is possible !!

 $\langle W \rangle_{\mathcal{T}} - \langle \Delta F \rangle_{\mathcal{T}} \geq -k_B T \langle \delta \rangle_{\mathcal{T}}$ if $\langle \delta \rangle_{\mathcal{T}} > 0$ (time-reversal asymmetry)

