

L2, Tuesday

Optimal Control

- If system is controllable, can put poles anywhere.
n-dim sys. \Rightarrow n poles \rightarrow where to put them?
- Rule of thumb: stabilize as needed; speed up slowest modes; leave rest untouched
- Here: define a "performance index" or "cost funct."
- choose control to minimize the cost

1d example:

$$J(u_0^\infty, x(0)) \equiv \int_0^\infty dt L(x, u) \equiv \int_0^\infty dt \cdot \frac{1}{2} (Qx^2(t) + Ru^2(t))$$

$$\dot{x} = -ax + u(t), \quad x(0) = x_0, \quad a > 0$$

- J is a functional of $u_0^\infty \equiv u(t), t \in [0, \infty)$
and a function of $x(0) = x_0$.
- need $J \geq 0$ (finite lower bound) $\Rightarrow Q \geq 0, R \geq 0$
 - $R=0 \Rightarrow$ want best performance (min $x(t)$)
regardless of control "effort"
 - $Q=0 \Rightarrow$ only care about cost (do nothing)

scale $Q \rightarrow 1$

Let $u = -Kx$ (can justify form...)

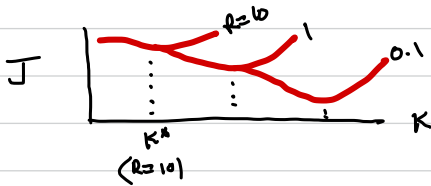
$$\Rightarrow \dot{x} = -(a+K)x \Rightarrow x(t) = x_0 e^{-(a+K)t} \Rightarrow J = x_0^2 \frac{(1+RK^2)}{4(a+K)}$$

$$\frac{\partial J}{\partial K} = 0 \Rightarrow K^2 + 2aK - \frac{1}{R} = 0 \Rightarrow K^* = -a \pm \sqrt{a^2 + \frac{1}{R}}$$

choose $+$ as $K^* > 0$

$Ra^2 \ll 1$: cheap control, $J \approx \int_0^{\infty} dt \frac{1}{2} \dot{x}^2$, $K^* \sim R^{-1/2} \rightarrow \infty$

$Ra^2 \gg 1$: expensive control, $J \approx \int_0^{\infty} dt \frac{1}{2} u^2$, $K^* \sim \frac{1}{2Ra} \rightarrow 0$



- Rather than choose K , choose R (or R/Q)
 - Have more intuition about R than K (here, a bit artificial)
 - Optimal \neq good! (poor $R \Rightarrow$ poor control)
- Q : unstable eigen?

General setting

$$J = \underbrace{\varphi[x(\tau)]}_{\text{terminal cost}} + \underbrace{\int_0^{\tau} dt L(x, u)}_{\text{running cost}}$$

• Dynamics $\dot{x} = f(x, u)$ as a constraint

• $J = J[u_0^{\tau}, x_0]$

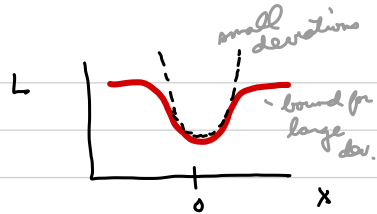
• $\tau \rightarrow \infty \Rightarrow$ drop $\varphi(\cdot)$

• soft constraint $\varphi(\cdot)$ vs. hard constraint BC

$$x(\tau) = x_{\tau}$$

\nearrow ag

- Running cost $L(x, u)$
 - scalar
 - bounded from below
 - smooth (eg., twice differentiable)



- Large φ , small L : "ends justify the means".

Impose constraints by Lagrange multipliers

— transpose of n -dim vector

$$J' \equiv \varphi[x(\tau)] + \int_0^\tau dt L(x, u) + \lambda^T(t) [f(x, u) - \dot{x}]$$

- Lagrange mult. $\lambda(t)$ is adjoint vector (costate)
- $\lambda(t)$ is a function of time because $\dot{x} = f$ must be imposed at every moment in time
- Solve unconstrained problem by allowing functional variations $\delta x(t)$, $\delta u(t)$, $\delta \lambda(t)$

$$\delta J' = (\partial_x \varphi) \delta x(\tau) + \int_0^\tau dt \left[(\partial_x L) \delta x + (\partial_u L) \delta u + \lambda^T \left((\partial_x f) \delta x + (\partial_u f) \delta u - \delta \dot{x} \right) + \delta \lambda^T (f - \dot{x}) \right]$$

$$= (\partial_x \varphi) \delta x(\tau) - \lambda^T(\tau) \delta x(\tau) + \lambda^T(0) \delta x(0) + \int_0^\tau dt \left[(\partial_x L) + \lambda^T (\partial_x f) + \dot{\lambda}^T \right] \delta x(t) + \int_0^\tau dt \left[(\partial_u L) + \lambda^T (\partial_u f) \right] \delta u(t) + \int_0^\tau dt (f - \dot{x})^T \delta \lambda(t)$$

use $a^T b = b^T a$

- The extra boundary terms come from integrating $\lambda^T(t) \delta \dot{x}(t)$ by parts
- $x(0) = x_0$ is fixed $\Rightarrow \delta x(0) = 0$
- In classical mechanics (CM), $u(t)$ is fixed $\Rightarrow x(\tau)$ fixed $\Rightarrow \delta x(\tau) = 0$ **not here!**

$\delta J' = 0 \Rightarrow$ Euler-Lagrange eqs.

- Use Thm: $\int_0^\tau dt f(t) \delta \varphi(t) = 0 \quad \forall \delta \varphi(t) \Rightarrow f(t) = 0 \quad (0 < t < \tau)$

$$0 \rightarrow \tau: \quad \dot{x} = f(x, u) \quad x(0) = x_0$$

$$\tau \rightarrow 0: \quad \dot{\lambda} = -(\partial_x f)^T \lambda - (\partial_x L)^T, \quad \lambda(\tau) = (\partial_x \phi|_\tau)^T$$

$$t: \quad 0 = (\partial_u f)^T \lambda + (\partial_u L)^T$$

- These eqs. give only **necessary** conditions 

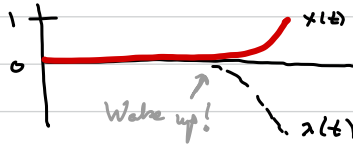
- Interpretation of adjoint $\lambda(t)$ (1d)
 - ignores λ -dependence in $(\partial_x f) \Rightarrow \dot{\lambda} = -(\partial_x f)\lambda - (\partial_x L)$
 - **Planning horizon** (timescale) $\sim -(\partial_x f)^T$
 - driven by $-\partial_x L$ term
 - **backwards in time**
- Together: Boundary-value problem (not init. val.)

1d pt. (again)

$$L = \frac{1}{2}(x^2 + u^2), \quad \dot{x} = -x + u$$

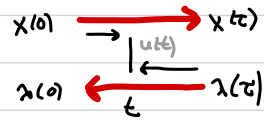
- but now let $x(0) = x_0$ and impose $x(\tau) = x_c$

Euler-Lagrange: $\dot{x} = -x + u, \quad \dot{\lambda} = +\lambda - x, \quad u = -\lambda$



• control signal $u(t)$

- depends on past states x_0^t
- depends on future plans λ^T



Ex: Swing up a pendulum

$$\ddot{\theta} + \sin\theta = u(t) \quad \text{torque control}$$

$$\begin{pmatrix} \theta \\ \dot{\theta} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \rightarrow \begin{pmatrix} \pi \\ 0 \end{pmatrix}$$

down up

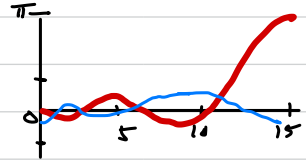
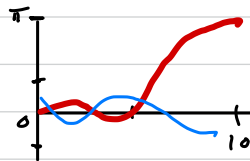
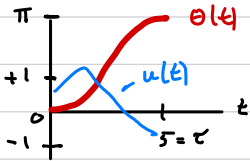
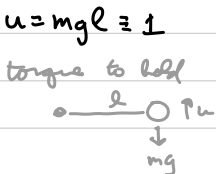
$$J = \int_0^\tau dt \cdot \frac{1}{2} u^2(t)$$

"control effort"

For DC motor, $u \sim \text{torque} \sim I$
if $I^*(t) = \text{torque}$

hard constraint on $x = \{\theta, \dot{\theta}\}$ at $t = \tau \Rightarrow$ no λ cond.

E-L: $\ddot{\lambda} + \lambda \cos\theta = 0, \quad u(t) = -\lambda(t) \rightarrow 4^{\text{th}}$ order BVP



- Increasing $\tau \Rightarrow$ reduce $|u|_{\max}$
- # reversals increases w/ τ

Hamiltonian formalism

Classical Mech:

Lagrangian $L(q, \dot{q}) = T - U$

kinetic en.
potential en.

Generalized momenta

$$p = \frac{\partial L}{\partial \dot{q}}$$

Legendre transf:

$$H = p\dot{q} - L$$

Hamilton's eqs:

$$\dot{q} = (\partial_p H)^T \quad \dot{p} = -(\partial_q H)^T$$

Control-theory version:

$$H = +L + \lambda^T \dot{x} = L + \lambda^T f$$

note change of sign

$$\dot{x} = (\partial_x H)^T = f$$

dynamics

$$\dot{\lambda} = -(\partial_x H)^T$$

adjoint

$$(\partial_u H)^T = 0$$

→ min H

- so adjoint states \leftrightarrow conjugate momenta
- simpler notation than Euler-Lagrange

Hamiltonian structure even if $\dot{x} = f(x, u)$ is dissipative

7.3 Linear Quadratic Regulator (LQR)

- small deviations, quadratic costs

$$J = \int_0^T dt \frac{1}{2} \left(\underbrace{x^T Q x}_L + \underbrace{u^T R u}_L \right)$$

Q, R symmetric
 $R > 0, Q \geq 0$
 add final cost if $Q \neq 0$

$$\dot{x} = f(x, u), \quad f(0, 0) = 0 \Rightarrow \dot{x} \approx Ax + Bu$$

$$H(x, \lambda, u) = \frac{1}{2} (x^T Q x + u^T R u) + \lambda (Ax + Bu)$$

$$\begin{aligned} \rightarrow \quad \dot{x} &= Ax + Bu, & \dot{\lambda} &= -A^T \lambda - Qx, & u &= -R^{-1} B^T \lambda \\ &(\partial_x H)^T & &(-\partial_x H)^T & \partial_u H = 0 \end{aligned}$$

• Can try

$$\lambda(t) = S(t) x(t)$$

n x n matrix
n-dim column vector
 this trick works only for linear dynamics!

$$u = -R^{-1} B^T S x \equiv -Kx$$

$$\dot{\lambda} = -Qx - A^T (\underbrace{Sx}_\lambda) = \frac{d}{dt}(Sx) = \dot{S}x + S(Ax - \underbrace{BR^{-1}B^T Sx}_x)$$

$$\dot{S} = -Q - A^T S - SA + SBR^{-1}B^T S$$

factor $x(t)$; holds $\forall t$
 continuous time Riccati Eq.

steady state

$$\rightarrow 0 = -Q - A^T S - SA + SBR^{-1}B^T S$$

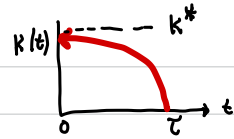
algebraic Riccati Eq.

• quadratic eq \Rightarrow multiple solns; only 1 is physical

Ex: Id control, yet again: $K = \bar{r}^{-1} (1) S = S/R$

$$\Rightarrow \dot{s} = -1 + 2a S + S^2/R$$

$$\rightarrow S = -aR \pm \sqrt{a^2 R^2 + R} \quad \text{steady state}$$



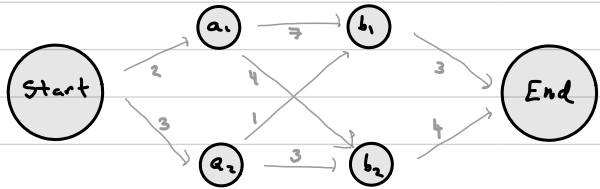
But fully optimal solution for finite τ has $K = K(t)$

$K \rightarrow 0$ as $t \rightarrow \tau$ (it costs for control, but no benefit if applied too close to end)

Dynamic Programming

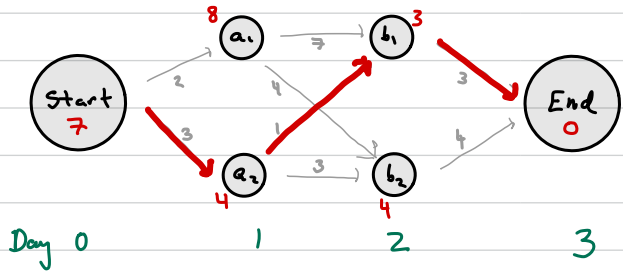
Richard Bellman, 1950s (RAND)

Shortest path
between cities



$$J(x) = \text{cost-to-go}$$

$$J^* = \text{optimal } J$$



$$J^*(\text{End}) = 0$$

You are already there.

$$J^*(b_1) = 3, \quad J^*(b_2) = 4$$

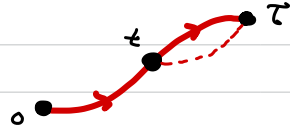
only one possibility in each case

$$J^*(a_1) = \min [7 + J^*(b_1), 4 + J^*(b_2)] = 8$$

$$J^*(a_2) = \min [1 + J^*(b_1), 3 + J^*(b_2)] = 4$$

$$J^*(\text{Start}) = \min [2 + J^*(a_1), 3 + J^*(a_2)] = 7$$

Principle of Optimality:



For any point on an optimal trajectory, the remaining trajectory is optimal, starting at that point.

Bellman Eq. (discrete case)

$X_k =$ state

city at time k ; eg a_1, a_2

$u_k =$ action

road choice at time k

$L(X_k, u_k) \equiv L_k$

running cost of current step, given X_k, u_k

$J(X_k, u_k) \equiv J_k$

cost-to-go; start from X_k , choose u_k, \dots, u_{N-1}

$\varphi(X_N)$

terminal cost

[note: $u_N = 0$]

$$J = \sum_{n=0}^{N-1} L(X_n, u_n) + \varphi(X_N) \quad \text{total cost}$$

$$\begin{aligned} J_k &= \sum_{n=k}^{N-1} L(X_n, u_n) + \varphi(X_N) \\ &= \underbrace{L(X_k, u_k)}_{L_k} + \underbrace{\sum_{n=k+1}^{N-1} L(X_n, u_n) + \varphi(X_N)}_{J_{k+1}} \end{aligned}$$

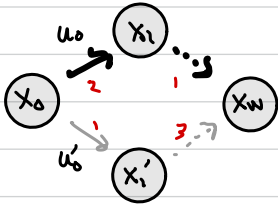
$$\begin{aligned} \text{So:} \quad J_k &= L_k + J_{k+1} & J_N &= \varphi(X_N) \\ X_{k+1} &= f(X_k, u_k) & X_0 &\text{ given} \end{aligned}$$

← J_k obeys a backwards recursion relation

- solving it fixes u_k^* (optimal choice at time k)
- then solve forward recursion reln for X_k →

Bellman eq.

$$J^*(x_k) = \min_{\{u_k\}} [L(x_k, u_k) + J^*(x_{k+1})]$$



We choose u_0 over u'_0 even though it has a higher immediate cost.

Overall future costs are lower

Solve J_k backwards; get u_k^* ; iterate $x_{k+1} = f(x_k, u_k^*)$ forwards.

- Solution is possible because dynamics have a state structure \Rightarrow solution decomposes into stages

Numerical algorithm: $\left. \begin{array}{l} \underline{n \text{ nodes / step}} \\ \Rightarrow \text{each node has } n \text{ possibilities} \end{array} \right\} \mathcal{O}(n^2)$
 $\times \underline{N \text{ stages}} \Rightarrow \mathcal{O}(N \cdot n^2)$

vs. naive search $\mathcal{O}(n^N)$ huge difference for large N !

\Rightarrow Value of planning.

"Life must be understood backwards but lived forwards."
 S. Kierkegaard

Bellman algorithm (re-) discovered many times

- Cell phones (Andrew Viterbi) Qualcomm. decoding
- Sequence alignment (bioinf.) (Needleman - Wunsch, Smith - Waterman)
- Shortest path (Dijkstra's alg.)
- \vdots

Bellman eq. (continuous case)

cost-to-go: $J(x, u) = \varphi[x(\tau)] + \int_t^\tau dt' L(x, u)$; $J^* = \inf_{u \in [t, \tau]} J(x, u)$

$$\begin{aligned} J^*(x) &= \inf_u \left[\varphi[x(\tau)] + \int_t^\tau dt' L(x, u) \right] \\ &= \inf_u \left[\int_t^{t+\Delta t} dt' L(x, u) + \varphi[x(\tau)] + \int_{t+\Delta t}^\tau dt' L(x, u) \right] \\ &\approx \inf_{u \in [t, t+\Delta t]} \left[L(x, u) \Delta t + \inf_{u \in [t+\Delta t, \tau]} [J(x + \dot{x} \Delta t, u(t+\Delta t))] \right] \\ &= \inf_{u \in [t, t+\Delta t]} \left[L(x, u) \Delta t + J^*(x + f \Delta t) \right] \\ &\approx \inf_{u \in [t, t+\Delta t]} \left[L(x, u) \Delta t + J^*(x) + (\partial_t J^*) \Delta t + (\partial_x J^*) f \Delta t \right] \end{aligned}$$

$$\partial_t J^* + \inf_{u(t)} \left[L(x, u) + (\partial_x J^*) f(x, u) \right] = 0$$

Hamilton - Jacobi - Bellman (HJB) eq.

- Integrate backwards in time from final condition
 $J^*[x(\tau), \tau] = \varphi[x(\tau)]$
 $\Rightarrow u^*(t)$

- Then integrate $\dot{x} = f(x, u^*)$ forward from $x(0) = x_0$

- $\partial_x J^* = n$ -dim row vector $\Rightarrow (\partial_x J^*) f$ is a scalar

One - dim control, yet again!

$$J = \int_t^{\tau} dt' \frac{1}{2}(x^2 + u^2), \quad \dot{x} = -x + u$$

$$\text{HJB:} \quad \partial_t J^* = - \inf_u \left[\underbrace{\frac{1}{2}(x^2 + u^2)}_{L(x,u)} + (\partial_x J^*) \underbrace{(-x + u)}_{f(x,u)} \right]$$

Since there are no restrictions on $u(t)$, we can do the minimization by taking ∂_u of right-hand side.

$$\begin{aligned} \partial_u \left[\frac{1}{2}(x^2 + u^2) + (\partial_x J^*)(-x + u) \right] &= u + (\partial_x J^*) = 0 \\ \Rightarrow u^* &= -\partial_x J^* \end{aligned}$$

$$\Rightarrow \partial_t J^* = -\frac{1}{2}x^2 + \frac{1}{2}(\partial_x J^*)^2 + x(\partial_x J^*)$$

Looks intimidating! But try $J^*(x, t) = \frac{1}{2}x^2 s(t)$

$$\Rightarrow \dot{s} = -1 + s^2 + 2s, \quad s(\tau) = 0$$

$$\Rightarrow u^* = -\partial_x J^* = -s x \equiv -kx$$

Now we can find $\dot{x} = -x - kx = -(1+k)x$, $x(0) = x_0$, etc.

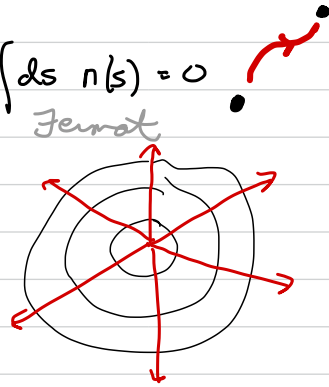
HJB vs. Euler-Lagrange

- Why do we have two apparently different ways to solve the optimal control problem?
- Same story in classical mechanics, optics, ...

Eg. , Optics

① Calculus of variations: $\delta \int ds n(s) = 0$ Fermat

② Eikonal (Huygen's wavefronts)
PDE for wavefronts.



So we can solve for rays (trajectories)
or wavefronts (cost functions)

Since rays are \perp wavefronts, methods are equiv.

- HJB usually much harder to solve, but you get sol'n for all x_0 at once.
- Euler-Lagrange is simpler but has to be redone for each x_0 .

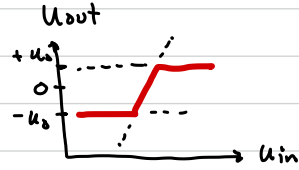
Hard constraints

soft: $J = \varphi[x(\tau)] + \dots$
 $\alpha = \dots \int \frac{1}{2} dt R u^2(t)$

VS.

hard: $x(\tau) = x_\tau$
 $|u(t)| \leq U_0$

hard constraint $|u| \leq U_0$ is a nonlinearity.



Strategies

Method 0: Buy a bigger actuator \rightarrow bigger u range

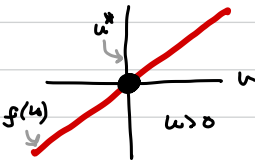
Pontryagin Minimum Principle: $\mathcal{H}(x, \lambda) = \inf_{u \in U} H(x, \lambda, u)$ \rightarrow dyn. programming

$$\dot{x} = (\partial_x \mathcal{H})^T, \quad \dot{\lambda} = -(\partial_x \mathcal{H})^T$$

$$x(0) = x_0, \quad \lambda(\tau) = (\partial_x \varphi)^T|_{t=\tau}$$

If $u(t)$ is not constrained $\inf_u \Rightarrow \partial_u = 0$

With constraint, $u(t)$ might get stuck at boundary



$$\inf_{u > 0} f(u) = f(0)$$

$$\frac{d}{dt} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} + \begin{pmatrix} 0 \\ 1 \end{pmatrix} u = \begin{pmatrix} \dot{x}_2 \\ u \end{pmatrix}$$

Ex. Moving in minimal time

- free particle in space: $\dot{x}_1 = x_2, \dot{x}_2 = u$ ($u = F/m$)

Goal: go from $\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_0 \rightarrow \begin{pmatrix} 0 \\ 0 \end{pmatrix}$ as fast as possible

Limits $|u(t)| \leq 1$ $J = \int_0^\tau dt (1)$ τ not fixed

Running cost $L=1 \Rightarrow H = 1 + \lambda_1 x_2 + \lambda_2 u \rightarrow (\lambda_1, \lambda_2) \begin{pmatrix} x_2 \\ u \end{pmatrix}$

$$\Rightarrow \dot{\lambda}_1 = -\frac{\partial H}{\partial x_1} = 0, \quad \dot{\lambda}_2 = -\frac{\partial H}{\partial x_2} = -\lambda_1$$

$$\Rightarrow \lambda_1 = \text{const}, \quad \lambda_2 = -c_1 t + c_2$$

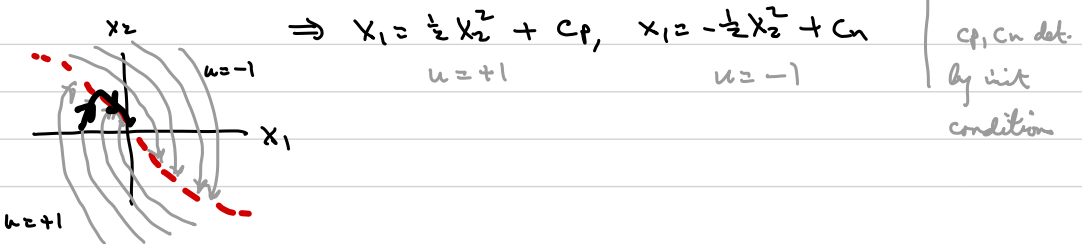
Can't min. H by $\frac{\partial H}{\partial u} = 0$, but can min. at boundary.

$\min_u H \Rightarrow u = -\text{sign}(\lambda_2)$ at all times

$\Rightarrow u(t)$ always $\geq \pm 1$

Since $\lambda(t)$ varies linearly, $u(t)$ switches once.

Solve eqs. of motion separately for $u = \pm 1$



Bang-bang control

In the previous problem, the control "lived" at $u = \pm u_0$

This behavior, "bang-bang control," results when $H(x, \lambda, u)$ is linear in u (or has no min. in range)

- often a bad idea in practice
 - hard on equipment
 - not robust [braking a car...]

Feedback

The calc. of variations approach (or Hamiltonian)

→ control $u^*(t)$ to take sys from $x_0 \rightarrow x_T$.

This is open loop, feedforward control.

via optimal $x(t)$

disturbances or modeling inaccuracies

→ trajectory will deviate from optimal

→ need feedback

- 1) Local linear feedback
- 2) Model Predictive Control

① Local linear feedback

$$J = \varphi[x(\tau)] + \int_0^{\tau} dt L(x, u), \quad \dot{x} = f(x, u), \quad x(0) = x_0$$

$$H = L + \lambda f, \quad \dot{\lambda} = -\partial_x H, \quad \partial_u H = 0$$

- Assume we have solved the "nominal feedforward" problem:
 $\Rightarrow x_{ff}(t), u_{ff}(t), \lambda_{ff}(t)$
- Perturb about nominal solution:

$$x(t) = x_{ff}(t) + \delta x(t), \quad u(t) = u_{ff}(t) + \delta u(t), \quad \lambda(t) = \lambda_{ff}(t) + \delta \lambda(t)$$

- Because the nominal "ff" solution is optimal, $\delta J = 0$

$$\Rightarrow \delta^2 J = \frac{1}{2} \delta x^T (\partial_{xx} \varphi) \delta x \Big|_{t=\tau} + \int_0^{\tau} dt (\delta x^T \delta u^T) \begin{pmatrix} Q & M \\ M & R \end{pmatrix} \begin{pmatrix} \delta x \\ \delta u \end{pmatrix}$$

with $Q(t) = \partial_{xx} L$, $M(t) = \partial_{xu} L$, $R(t) = \partial_{uu} L$
 - evaluated along $x_{ff}(t)$ and $u_{ff}(t)$

Augment the cost pert. w/ perturbed dynamics. ie $Q = \frac{\partial^2 L(x, u)}{\partial x^2} \Big|_{x_{ff}, u_{ff}}$

$$\delta^2 J' = \delta^2 J + \delta \lambda^T \left[A \delta x + B \delta u - \dot{\delta x} \right] \quad A(t) = \partial_x f, \quad B(t) = \partial_u f$$

$$\text{E-L: } \delta \dot{\lambda} = -Q \delta x - M \delta u - A^T \delta \lambda, \quad \delta u = -R^{-1} (M \delta x + B^T \delta \lambda)$$

$-\partial_x H$ $-\partial_u H = 0$

This is just a slightly more complicated LQR problem.

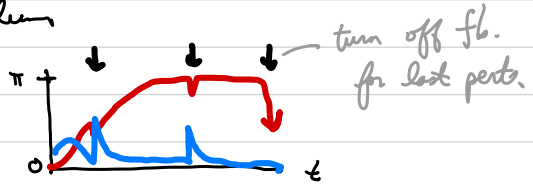
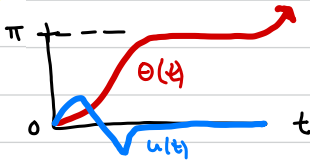
- all matrices time dependent $A \rightarrow A(t)$, etc.
- quadratic costs have cross-term $M(t)$ ($\delta x^T \cdot M \cdot \delta x$)

Fortunately, all steps of LQR carry through analogously.

full control signal is then

$$u(t) = u_{ff}(t) + u_{fb}(t) = u_{ff}(t) + K(t) [x_{ff}(t) - x(t)]$$

Ex: Swing up + balance pendulum



Model Predictive Control (MPC)

- Compute ff from $x(t)$ to $x(t+\tau) \Rightarrow u_{ff}(t)$
- Apply $u_{ff}(t)$ for short time ("one step")

"Feedback by repeated ff"

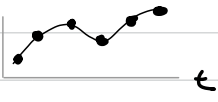
- good for problems w/ constraints
- costly to compute \Rightarrow good for "slow" problems
- widely used in "slow" industries (eg chemical plants)

Numerical Methods

→ direct: $\min J[u]$
 vs. indirect: $\delta J = 0$



Direct: $\min_{u(t) \in U} J[u(t); x_0]$

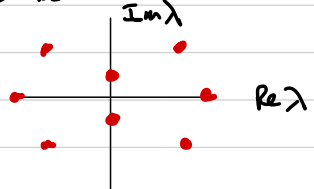
Project $u(t)$ on finite set $u(t)$ 

n : $u(t) = \sum_n a_n \varphi_n(t)$ basis funct.

Then solve directly (often non-convex )

Indirect: solve rec. eqs. (Hamilton, PMP)

$n \rightarrow 2n + n$ constraints dim

Hamiltonian H real \Rightarrow 

\Rightarrow stiff eqs.

(beard for shooting methods)

\rightarrow solve discretization by Newton's method

\rightarrow can write Jacobian as band diag. $\Rightarrow O(N)$
 can even solve each stage in \parallel !